# The Takesaki equivalence relation for maximal abelian subalgebras

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**Abstract.** For a maximal abelian subalgebra  $A \subset M$  in a finite von Neumann algebra, we consider an invariant due to Takesaki which is an equivalence relation on a standard probability space. We give several characterizations of this invariant and show that it can be reconstructed from the A-bimodule structure of the GNS Hilbert space  $L^2(M)$ . In particular, we show that this invariant is induced by the action of the normalizer on A. Hence, this gives a new proof to a question of Takesaki.

#### 1. Introduction

In this paper, we study maximal abelian subalgebras (MASAs) in a finite von Neumann algebra with a separable predual. We will always denote such an inclusion as  $A \subset M$  and we fix a faithful normal unital trace  $\tau$  on M. The study of MASAs began with the work of Dixmier in 1954 [3]. (See [14] for a general introduction on the subject.) He considered the normalizer  $N_M(A)$  which is the group of unitaries  $u \in M$  such that  $uAu^* = A$ . In 1963, Takesaki introduced in [15] a measure theoretical invariant for a MASA. An explicit presentation of this invariant is given below.

Let us first define the Takesaki equivalence relation. We refer the reader to [16, Chap. 4, § 8] for a presentation of the general theory of direct integrals of Hilbert spaces, representations, and von Neumann algebras. Let Y be a compact Hausdorff space and  $\nu$  a Borel probability measure on it such that A is isomorphic to the von Neumann algebra  $L^{\infty}(Y,\nu)$ . We fix such an isomorphism and identify the two von Neumann algebras. Let  $L^2(M)$  be the GNS Hilbert space associated to the trace  $\tau$  and  $x \mapsto x\Omega$  the embedding of M in  $L^2(M)$ . Let  $\pi$ ,  $\rho$  be the left and right actions of M on the Hilbert space  $L^2(M)$ , i.e.  $\pi(x)\rho(y)(z\Omega) = xzy\Omega$ . Consider the measurable field of Hilbert

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spaces  $\{\mathcal{K}_t, t \in Y\}$  such that  $L^2(M)$  is equal to the direct integral

$$\int_{Y}^{\oplus} \mathcal{K}_t d\nu(t),$$

such that  $\rho(A)$  becomes the algebra of all diagonalizable operators. Let  $B \subset M$  be a separable  $C^*$ -subalgebra that is dense for the weak topology. Consider the measurable field of representations of B,  $\{\pi_t, t \in Y\}$ , such that

$$\pi|_{B} = \int_{Y}^{\oplus} \pi_{t} d\nu(t),$$

where  $\pi|_B$  denotes the restriction to B of the standard representation.

**Definition 1.1.** Let  $\mathcal{R}$  be the equivalence relation on Y defined by  $(s,t) \in \mathcal{R}$  if and only if the representation  $\pi_s$  is unitarily equivalent to  $\pi_t$ . It is the *Takesaki* equivalence relation.

We write  $\pi_s \simeq \pi_t$  to say that the two representations are unitarily equivalent.

**Definition 1.2.** Let  $E, F \subset Y^2$  be some subsets, we say that E is weakly contained in F if there exists a null set  $N \subset Y$  such that  $E \setminus N^2 \subset F$ , where  $E \setminus N^2 = \{x \in E, x \notin N^2\}$ . We denote this by  $E \prec F$ . This defines a partial order. We say that E is equivalent to F if  $E \prec F$  and  $F \prec E$  and denote it by  $E \equiv F$ . This defines an equivalence relation on the subsets of  $Y^2$ . We denote the equivalence class of a subset E by  $\widehat{E}$ .

**Definition 1.3.** Let  $\widehat{\mathcal{R}}$  be the equivalence class of  $\mathcal{R}$  for  $\equiv$ . It is an invariant for the MASA  $A \subset M$  that we call the *Takesaki invariant*. In particular, it does not depend of the choice of the  $C^*$ -algebra B, see [15]. We say that a MASA is *Takesaki simple* if  $\mathcal{R} \equiv \Delta Y$ , where  $\Delta Y = \{(t,t), t \in Y\}$  is the diagonal of Y.

Let us define an other equivalence relation. Consider the normalizer  $N_M(A)$  and a countable subgroup  $G < N_M(A)$  such that the bicommutant  $\{G \cup A\}''$  is equal to  $N_M(A)''$  inside M. The group G acts on A, hence this gives an action on the space  $(Y, \nu)$ . We denote by  $\mathcal{N}_G$  the orbit equivalence relation. This equivalence relation does not depend of the choice of the group G (see Proposition 2.2), therefore we simply denote it by  $\mathcal{N}$ .

Takesaki proved in [15, Thm. 1.2] that  $\mathcal{N} \prec \mathcal{R}$ . He asked if  $\mathcal{R}$  is a countable, quasi-invariant equivalence relation and if  $\mathcal{R} \equiv \mathcal{N}$ . We recall that an equivalence relation is quasi-invariant if the saturation of a null set is still a null set. In the mid 70's, Hahn developed a theory of measure groupoids [8], [9]. Using this theory, he proved in [10] that  $\mathcal{R} \equiv \mathcal{N}$ .

The author wants to indicate that all the results presented in this paper have been proved without knowing the work of Hahn. We give here an elementary proof of the equivalence  $\mathcal{R} \equiv \mathcal{N}$ . Furthermore, we define a third equivalence relation  $\mathcal{B}$  that we call the *bimodule equivalence relation* and prove that  $\mathcal{R} \equiv \mathcal{B}$ . In particular, we show the surprising fact that the Takesaki equivalence relation can be reconstructed from the A-bimodule structure of  $L^2(M)$ .

The bimodule equivalence relation is defined in a similar way as  $\mathcal{R}$  except that we replace the separable weakly dense  $C^*$ -subalgebra  $B \subset M$  by a separable weakly dense  $C^*$ -subalgebra  $D \subset A$  (see Definition 2.1).

One of the key arguments to prove this theorem is to show that the Takesaki equivalence relation is quasi-invariant. We introduce a subset  $\mathcal{Y} \subset Y \times Y$  that we call the set of atoms (see Definition 2.4). This set has been studied in [13] and in [12]. We identify  $\mathcal{Y}$  as a symmetric relation on Y and prove that it is quasi-invariant. Furthermore, we show that  $\mathcal{B}$  is weakly contained in  $\mathcal{Y}$ . Therefore,  $\mathcal{R}$  and  $\mathcal{B}$  are quasi-invariant. We also prove that  $\mathcal{Y}$  is equivalent to  $\mathcal{R}$  under the relation " $\equiv$ ", hence the main theorem of this paper is the following.

**Theorem 1.4.** Let  $A \subset M$  be a MASA in a finite von Neumann algebra, then the Takesaki equivalence relation, the bimodule equivalence relation, the set of atoms and the equivalence relation induced by the normalizer are equal up to a null set, i.e.

$$\mathcal{R} \equiv \mathcal{B} \equiv \mathcal{V} \equiv \mathcal{N}$$
.

In particular, a MASA is singular (its normalizer is equal to the unitaries of A) if and only if it is Takesaki simple. We illustrate this result with a proposition on inclusions of countable groups, see Proposition 4.1. Furthermore, we deduce a result on the normalizer of tensor product of MASAs.

Corollary 1.5. Consider a family of MASAs in some von Neumann algebras  $\{A_l \subset M_l, l \in \Lambda\}$ , where  $\Lambda$  is a countable set. For any  $l \in \Lambda$ , we consider a faithful trace  $\tau_l$  on  $M_l$ . Let  $A = \bigotimes_l A_l$  and  $M = \bigotimes_l M_l$  be the tensor products of those von Neumann algebras with respect to the traces  $\tau_l$ . The von Neumann algebra generated by the normalizer  $N_M(A)$  is equal to the tensor product of the von Neumann algebras generated by the normalizers  $N_{M_l}(A_l)$ , i.e.

$$N_M(A)^{\prime\prime} = \bigotimes_{l \in \Lambda} N_{M_l}(A_l)^{\prime\prime}.$$

This result has been proved by Chifan [2] using analytic techniques.

The rest of the paper is organized into 3 sections. In Section 2 we fix some notations and review some basic facts about equivalence relations. We define the equivalence relation induced by the action of the normalizer on the algebra A, the bimodule equivalence relation and the set of atoms. In Section 3 we prove the main result of this paper. In Section 4 we illustrate the main result in the context of representations of discrete countable groups. Then we prove the corollary on tensor product of MASAs.

#### 2. Notations and definitions

In this section, we fix some notations and define the bimodule equivalence relation  $\mathcal{B}$ , the equivalence relation  $\mathcal{N}$  induced by the action of the normalizer and the set of atoms  $\mathcal{Y}$ . Consider the  $C^*$ -algebra of continuous functions  $\mathcal{C}(Y)$ . Let  $B \subset M$  be a  $C^*$ -subalgebra which is separable and weakly dense.

The equivalence class of  $\mathcal{R}$  does not depend on B. Therefore we can assume that B contains the  $C^*$ -algebra  $\mathcal{C}(Y)$ . Hence we have the following square of inclusions.

$$\begin{array}{ccc} L^{\infty}(Y,\nu) & \subset & M \\ & \cup & & \cup \\ \mathcal{C}(Y) & \subset & B \end{array}$$

**Definition 2.1.** Consider the bimodule equivalence relation  $\mathcal{B}$  which is defined such that  $(s,t) \in \mathcal{B}$  if the representation  $\pi_t|_{\mathcal{C}(Y)}$  is unitarily equivalent to  $\pi_s|_{\mathcal{C}(Y)}$ .

Let  $\mathcal{I}(Y,\nu)$  be the group of Borel automorphisms of  $(Y,\nu)$  that preserve the class of the measure  $\nu$ . For any countable subgroup  $H < N_M(A)$  we can define a group homomorphism  $\Theta^H : H \longrightarrow \mathcal{I}(Y,\nu)$  such that for any  $u \in H$  and any  $f \in A$ ,  $u^*fu = f \circ \Theta^H_u \nu$ -almost everywhere  $(\nu$ -a.e.). Consider the orbit equivalence relation

$$\mathcal{N}_H = \{ (\Theta_u^H(t), t), \ t \in Y, \ u \in H \}.$$

If  $u \in N_M(A)$  we denote by  $\Theta_u$  a given automorphism such that  $u^*fu = f \circ \Theta_u \ \nu$ -a.e. The following proposition justifies the definition of  $\mathcal{N}$  given in the introduction.

**Proposition 2.2.** There exists a countable subgroup  $G < N_M(A)$  such that  $G'' = N_M(A)'' \subset M$ . We denote by  $\mathcal{N}$  the orbit equivalence relation  $\mathcal{N}_G$ . If  $H < N_M(A)$  is a countable subgroup then  $\mathcal{N}_H \prec \mathcal{N}$ . Furthermore,  $\mathcal{N}_H \equiv \mathcal{N}$  if and only if

$$\{H \cup A\}'' = N_M(A)''.$$

Before proving this proposition we recall a useful lemma.

**Lemma 2.3.** Consider a unitary  $u \in N_M(A)$ . Suppose that there exists a Borel subset  $E \subset Y$  such that for any  $t \in E$ ,  $\Theta_u(t) \neq t$ . Then,  $\tau(u\chi_E) = 0$ , where  $\chi_E$  is the characteristic function of the set E.

*Proof.* Let us show that  $E_A(u\chi_E) = 0$ . Let  $f \in L^{\infty}(Y,\nu)$  be an injective function. We have that

$$fE_A(u\chi_E) = E_A(fu\chi_E) = E_A(u(u^*fu)\chi_E) = E_A(u(f \circ \Theta_u)\chi_E)$$
  
=  $E_A(u\chi_E)(f \circ \Theta_u) = (f \circ \Theta_u)E_A(u\chi_E).$ 

We identify  $E_A(u\chi_E)$  with a function of the algebra  $L^{\infty}(Y,\nu)$ . We have that  $(f-f\circ\Theta_u)(t)E_A(u\chi_E)(t)=0$  a.e. The function f is injective and  $\Theta_u(t)\neq t$  for any  $t\in E$ . Therefore,  $E_A(u\chi_E)(t)=0$  a.e., hence  $E_A(u\chi_E)=0$ . This implies that  $\tau(u\chi_E)=\tau\circ E_A(u\chi_E)=0$ .

Let us prove the proposition.

Proof of Proposition 2.2. Let  $G < N_M(A)$  be a countable subgroup which is dense for the norm of  $L^2(M)$ . This group satisfies that  $G'' = N_M(A)''$ . To

prove the two other statements of the proposition it is sufficient to show that for any countable subgroups  $H, K < N_M(A)$  we have that

$$\{H \cup A\}'' \subset \{K \cup A\}''$$

if and only if  $\mathcal{N}_H \prec \mathcal{N}_K$ .

Suppose that  $\mathcal{N}_H \prec \mathcal{N}_K$ , let  $u \in \{H \cup A\}''$  be a unitary. Let us show that  $u \in \{K \cup A\}''$ . Let  $\{v_k, k > 0\}$  be an enumeration of the countable group K. Consider the sets

$$E_k = \{ t \in Y, \, \Theta_{v_k}^K(t) = \Theta_u(t) \}$$

and

$$F_k = E_k \setminus \bigcup_{j < k} E_j.$$

We see immediately that the sets  $F_k$  are measurable. Let  $p_k$  be the projection equal to the characteristic function of the set  $F_k$ . If  $k \neq l$ , then  $v_k p_k \perp v_l p_l$ . Therefore, the sum

$$\sum_{k=1}^{\infty} v_k p_k$$

converges in the von Neumann algebra  $\{K \cup A\}''$  to an element v. By hypothesis, the graph of  $\Theta_u$  is weakly contained in  $\mathcal{N}_K$ . This implies that  $\mathcal{Y} \setminus \bigcup_k F_k$  is a null set. Thus,  $v = \sum_k v_k p_k$  is a unitary in  $\{K \cup A\}''$  and by construction  $\Theta_u = \Theta_v \ \nu$ -a.e. Therefore,  $vu^*$  is a unitary of A, thus  $u \in \{K \cup A\}''$ .

Suppose that  $\{H \cup A\}'' \subset \{K \cup A\}''$ . Consider a unitary  $u \in H$  and the set

$$E = \{ t \in Y, (\Theta_u^H(t), t) \notin \mathcal{N}_K \}.$$

The set E is measurable. Let v=up, where  $p=\chi_E$ . Consider a unitary  $w\in K\cup U(A)$ , where U(A) is the unitary group of A. Let us show that v is orthogonal to w, i.e.  $\tau(w^*v)=0$ . We have that  $\Theta_{w^*u}=\Theta_w^{-1}\circ\Theta_u^H$  a.e. By assumption, for any  $t\in E$ ,  $\Theta_u^H(t)\neq\Theta_w(t)$  a.e., hence  $\Theta_{w^*u}(t)\neq t$  a.e. We can apply Lemma 2.3; thus,  $\tau(w^*v)=0$ . Therefore, the partial isometry v is orthogonal to the von Neumann algebra  $\{K\cup A\}''$ . This implies that p=0 and so the graph of  $\Theta_u^H$  is weakly contained in the equivalence relation  $\mathcal{N}_K$ . So,  $\mathcal{N}_H \prec \mathcal{N}_K$ .

We define the set of atoms. Let  $\mathcal{A} = \{\pi(A) \cup \rho(A)\}'' \subset \mathbb{B}(L^2(M))$  be the abelian von Neumann subalgebra generated by the left and right actions of A on  $L^2(M)$ . Consider the coordinate projection  $p: Y^2 \longrightarrow Y$ , p(s,t) = t and the flip  $\theta: Y^2 \longrightarrow Y^2$ ,  $\theta(s,t) = (t,s)$ . Following the proof of [5, Thm. 1], there exists a Borel probability measure  $\mu$  on  $Y^2$  such that the von Neumann algebra  $\mathcal{A}$  is isomorphic to  $L^{\infty}(Y^2,\mu)$ . It is easy to see that  $\mu$  is quasi-invariant with respect to the flip, i.e.  $\theta_*\mu \approx \mu$ . Furthermore, the push-forward measure  $p_*\mu$  is in the equivalence class of the measure  $\nu$ , therefore by [1, Chap. 6, § 3] there exists a disintegration of  $\mu$  with respect to  $(p,\nu)$ . It means that there exists a

unique a.e. family  $\{\mu_t, t \in Y\}$  of probability measures on Y, such that for any positive measurable function  $f: Y^2 \longrightarrow \mathbb{R}_+$ , the map

$$t \mapsto \int_Y f(s,t) d\mu_t(s)$$

is measurable, and

$$\mu(f) = \int_{Y} \int_{Y} f(s,t) d\mu_t(s) d\nu(t).$$

**Definition 2.4.** The set of atoms of the MASA  $A \subset M$  is the set

$$\mathcal{Y} = \{(s,t) \in Y^2, \, \mu_t(\{s\}) > 0, \, \mu_s(\{t\}) > 0\}.$$

The set  $\mathcal{Y}$  defines a symmetric relation on Y, hence we call orbit of t the set of  $s \in Y$  such that  $(s,t) \in \mathcal{Y}$ . Note that  $\mathcal{Y}$  is a measurable subset of  $Y^2$ , see [12, Prop. 3.3] for a proof.

### 3. The main result

**Theorem 3.1.** Consider the equivalence relations  $\mathcal{R}$ ,  $\mathcal{B}$ ,  $\mathcal{N}$  and the set of atoms  $\mathcal{Y}$ . Then,

$$\mathcal{R} \equiv \mathcal{B} \equiv \mathcal{Y} \equiv \mathcal{N}.$$

*Proof.* By definition,  $\mathcal{R} \subset \mathcal{B}$ . Let us show that  $\mathcal{B} \prec \mathcal{Y}$ . Consider a continuous function  $f \in \mathcal{C}(Y)$ .

**Claim:** The scalar f(t) is an eigenvalue of the operator  $\pi_t(f)$   $\nu$ -a.e.

Proof of the claim: The inclusion  $A \subset M$  gives us an inclusion of A-bimodules  ${}_{A}L^{\infty}(Y,\nu)_{A} \subset {}_{A}L^{2}(M)_{A}$ . As a right A-module,

$$L^{\infty}(Y,\nu)_A = \int_Y^{\oplus} \mathbb{C}_t d\nu(t),$$

where  $\mathbb{C}_t$  is the complex vector space of dimension one and  $\pi_t(f)$  acts by multiplication by f(t) on it  $\nu$ -a.e. Therefore, f(t) is an eigenvalue of  $\pi_t(f)$   $\nu$ -a.e.

The von Neumann algebra  $\mathcal{A}$  is isomorphic to  $L^{\infty}(Y^2, \mu)$  and it acts on the Hilbert space  $L^2(M)$ . Hence, there exists a measurable field of Hilbert spaces  $\mathcal{H}_{s,t}$  over  $(Y^2, \mu)$  such that we have an isomorphism of  $\mathcal{A}$ -modules

$$\phi: L^2(M) \simeq \int_{Y^2}^{\oplus} \mathcal{H}_{s,t} d\mu(s,t).$$

We define the following direct integral of Hilbert spaces

$$\mathcal{K}_t = \int_Y^{\oplus} \mathcal{H}_{s,t} d\mu_t(s).$$

By a result of Guichardet [7, Prop. 1], we have an isomorphism of right A-modules

$$L^2(M)_A \simeq \int_Y^{\oplus} \mathcal{K}_t d\nu(t).$$

Consider a continuous and injective function  $f \in \mathcal{C}(Y)$ . For any  $t \in Y$ , we associate to f an operator  $f_t \in \mathbb{B}(\mathcal{K}_t)$  determined by

$$f_t \xi = \int_Y^{\oplus} f(s) \xi_s d\mu_t(s),$$

for any

$$\xi = \int_{Y}^{\oplus} \xi_{s} d\mu_{t}(s) \in \int_{Y}^{\oplus} \mathcal{H}_{s,t} d\mu_{t}(s).$$

We remark that

$$\pi(f) = \int_{V}^{\oplus} f_t d\nu(t).$$

By uniqueness of the disintegration there exists a null set  $N \subset Y$  such that  $\pi_t(f) = f_t$  for any  $t \in Y \setminus N$ . The claim implies that there exists a null set  $N_0$  such that for any  $t \in Y \setminus N_0$  we have that f(t) is an eigenvalue of  $\pi_t(f)$ . Let  $(s_0, t) \in \mathcal{B} \setminus (N \cup N_0)^2$ , hence we have that  $f(s_0)$  is an eigenvalue of the operator  $f_t$ . Then there exists a non null vector

$$\eta = \int_{Y}^{\oplus} \eta_s d\mu_t(s) \in \mathcal{K}_t$$

such that  $f_t(\eta) = f(s_0)\eta$ , meaning that  $(f(s) - f(s_0))\eta_s = 0$   $\mu_t$ -a.e. This implies that  $s_0$  is an atom of  $\mu_t$  because f is injective. By exchanging the role of  $s_0$  and t we get that  $(s_0, t) \in \mathcal{Y}$ . Therefore,  $\mathcal{B} \prec \mathcal{Y}$ .

Let us show that  $\mathcal{Y} \prec \mathcal{N}$ .

**Claim:** Let  $X \subset \mathcal{Y}$  be a measurable subset. The following assertions are equivalent:

- (1) X is a null set for  $\mu$ ;
- (2)  $p_1(X)$  is a null set for  $\nu$ ;
- (3)  $p_2(X)$  is a null set for  $\nu$ .

Proof of the claim: we have that

$$\mu(X) = \int_Y \mu_t(\{s,\ (s,t) \in X\}) d\nu(t) = \int_{p_2(X)} \mu_t(\{s,\ (s,t) \in X\}) d\nu(t).$$

The set X is contained in  $\mathcal{Y}$ , hence for any  $t \in p_2(X)$ ,  $\mu_t(\{s, (s,t) \in X\})$  is strictly positive. Therefore  $\mu(X) = 0$  if and only if  $p_2(X)$  is a null set. We know that the class of the measure  $\mu$  is invariant under the flip; thus,  $\mu(X) = 0$  if and only if  $\mu(\theta(X)) = 0$  if and only if  $p_1(X)$  is a null set.

Let  $h: N \subset Y \longrightarrow Y$  be a measurable map defined on a measurable subset  $N \subset Y$  such that its graph  $\Gamma_h$  is weakly contained in  $\mathcal{Y}$ . We have that  $p_1(\Gamma_h) = h(N)$  and  $p_2(\Gamma_h) = N$ . Thus by the claim we have that  $\nu(N) = 0$  if and only if h(N) is a null set.

The set of atoms is measurable, hence there exists a  $\mu$ -null set  $N \subset Y^2$  such that  $\mathcal{Y} \setminus N$  is a Borel set. By [11, Thm. 18.10] there exists a countable family of Borel automorphisms  $h_k \in \mathcal{I}(Y, \nu)$  such that

$$\mathcal{Y}\setminus N=\bigcup_k\Gamma_{h_k}.$$

The claim implies that there exists a null set  $N_0 \subset Y$  such that  $N \subset N_0^2$ . Therefore  $\mathcal{Y}$  is equivalent to the union of graphs  $\bigcup_k \Gamma_{h_k}$ . To conclude that  $\mathcal{Y} \prec \mathcal{N}$ , we need to show that for any  $h_k$  there exists a unitary  $u_k \in N_M(A)$  such that  $\Theta_{u_k} = h_k \ \nu$ -a.e. Hence, by considering the subgroup of  $N_M(A)$  generated by the  $u_k$  we will get that  $\mathcal{Y} \prec \mathcal{N}$ .

**Lemma 3.2.** Consider a Borel automorphism  $h \in \mathcal{I}(Y, \nu)$  such that its graph  $\Gamma_h$  is weakly contained in  $\mathcal{Y}$ . Then, there exists a unitary in the normalizer of A,  $u \in N_M(A)$ , such that  $\Theta_u = h \nu$ -a.e.

*Proof.* Let us write  $L^2(M)$  as a direct integral of Hilbert spaces over the measure space  $(Y^2, \mu)$ 

$$\psi: L^2(M) \longrightarrow \int_{Y^2}^{\oplus} \mathcal{H}_{s,t} d\mu(s,t).$$

Consider a vector

$$\xi^0 = \int_{V^2}^{\oplus} \xi_{s,t}^0 d\mu(s,t)$$

such that  $\|\xi_{s,t}^0\|_{\mathcal{H}_{s,t}} = 1$   $\mu$ -a.e. Then the A-bimodule generated by the vector  $\xi^0$  is giving us an embedding of bimodules  $L^2(Y^2,\mu) \subset L^2(M)$ . Let  $\xi = \chi_{\Gamma_h} \in L^{\infty}(Y^2,\mu) \subset L^2(M)$  be the characteristic function of the graph of h. The preceding claim implies that  $\mu(\Gamma_h)$  is strictly positive. It is easy to see that  $\|\xi\|_2^2 = \mu(\Gamma_h)$ , hence  $\xi$  is different from zero. Furthermore, for any  $f \in A$ ,

$$f.\xi = \xi.f \circ h^{-1}.$$

The vector  $\xi$  is an affiliated operator to M, therefore it admits a spectral decomposition  $\xi = u|\xi|$ , where u is a partial isometry of M and  $|\xi| \in L^2(M)_+$ . Let p be the support of the partial isometry u. It is a projection of A, hence there exists a measurable subset  $E \subset Y$  such that  $1 - p = \chi_E$ . We have that

$$\|q.\xi\|^2 = \int_{V^2}^{\oplus} \chi_E(s) \chi_{\Gamma_h}(s,t) d\mu(s,t) = \mu(\{(s,h(s)), s \in E\}).$$

By the preceding claim,  $\mu(\{(s,h(s)), s \in E\}) = 0$  if and only if  $\nu(E) = 0$ . Therefore, 1-p=0. This implies that u is a unitary. We have that  $ufu^* = f \circ h$  for any  $f \in A$ .

Therefore, we have that  $\mathcal{Y} \prec \mathcal{N}$ .

Let us show that  $\mathcal{N} \prec \mathcal{R}$ . Let  $G < N_M(A)$  be a countable subgroup of  $N_M(A)$  that implements the orbit equivalence relation  $\mathcal{N}$ . Let  $u \in G$ , by [15, Thm. 1.2], the graph of  $\Theta_u$  is weakly contained in  $\mathcal{R}$ . Therefore,  $\mathcal{N} \prec \mathcal{R}$ . This achieves the proof of the theorem.

**Corollary 3.3.** A MASA in a finite von Neumann algebra is singular if and only if it is Takesaki simple.

*Proof.* This is a direct consequence of the fact that  $\mathcal{R} \equiv \mathcal{N}$ .

**Remark 3.4.** Consider a A-bimodule  ${}_{A}\mathcal{H}_{A}$ . We can define an equivalence relation as follows. Let  $\{\mathcal{K}_{t}, t \in Y\}$  be a measurable field of Hilbert spaces such that the right A-module  $\mathcal{H}_{A}$  is isomorphic to the direct integral

$$\int_{Y}^{\oplus} \mathcal{K}_t d\nu(t).$$

Let  $\lambda: A \longrightarrow \mathbb{B}(\mathcal{H})$  be the left action of A. Consider the measurable field of representations of  $\mathcal{C}(Y)$ ,  $\{\lambda_t, t \in Y\}$ , such that

$$\lambda = \int_{Y}^{\oplus} \lambda_t d\nu(t).$$

Let  $\mathcal{B}_{\mathcal{H}}$  be the equivalence relation  $\{(s,t) \in Y^2, \lambda_t \simeq \lambda_s\}$ . If  $\mathcal{H} = L^2(M)$  as A-bimodules, we get that  $\mathcal{B}_{\mathcal{H}}$  is the bimodule equivalence relation  $\mathcal{B}$ . Therefore, the bimodule equivalence relation can be reconstructed from  ${}_AL^2(M)_A$ , and so does the equivalence relation  $\mathcal{N}$  by Theorem 3.1. Feldman and Moore [4, 5] showed that a Cartan subalgebra is characterized by the equivalence relation  $\mathcal{N}$  and a 2-cocycle. Hence, a Cartan subalgebra is characterized by the A-bimodule  $L^2(M)$  and a 2-cocycle.

## 4. Illustration and consequences of the main theorem

4.1. Group von Neumann algebras. Let H < G be an inclusion of discrete countable groups such that H is abelian and

(2) for any 
$$g \in G \setminus H$$
 the set  $\{hgh^{-1}, h \in H\}$  is infinite.

This implies that the inclusion of group von Neumann algebras  $L(H) \subset L(G)$  is a MASA in a finite von Neumann algebra (see Godement [6]). Let  $C \subset G$  be a system of left coset representatives of H. We denote the unity of G by 1 and assume that  $1 \in C$ . Consider two functions  $\sigma: G \to C$  and  $\eta: G \to H$  such that for any  $g \in G$  we have  $g = \sigma(g)\eta(g)$ . If I is a set we denote the Hilbert space of square summable complex valued functions on I by  $\ell^2(I)$ . We denote the standard basis of  $\ell^2(G)$  (resp.  $\ell^2(H)$ , resp.  $\ell^2(C)$ ) by  $\{e_g, g \in G\}$  (resp.  $\{e_h, h \in H\}$  resp.  $\{\varepsilon_c, c \in C\}$ ). Let  $\pi, \rho: G \to \ell^2(G)$  be the left and right regular representations of the group G.

Let us decompose  $\pi$  with respect to the right action of H. Consider the dual group of H with its Haar measure  $(\widehat{H}, \nu)$  and the Fourier transform  $\mathcal{F}: \ell^2(H) \longrightarrow L^2(\widehat{H}, \nu)$  defined such that  $\mathcal{F}(e_h)(t) = t(h)$  for any  $h \in H$  and any character  $t \in \widehat{H}$ . We have an isomorphism of right L(H)-modules  $\phi: \ell^2(G) \longrightarrow \ell^2(C) \otimes L^2(\widehat{H}, \nu)$  given by the formula

$$\phi(e_g)(t) = t(\eta(g))\varepsilon_{\sigma(g)}.$$

Consider the representation  $\pi_t: G \curvearrowright \ell^2(C)$  defined such that

$$\pi_t(g)\varepsilon_c = t(\eta(gc))\varepsilon_{\sigma(gc)},$$

for any  $t \in \widehat{H}$ ,  $g \in G$  and  $c \in C$ . An easy computation shows that this gives a disintegration of the representation  $\pi$  with respect to the right action of H, i.e.

 $\pi = \int_{\widehat{H}} \pi_t d\nu(t).$ 

**Proposition 4.1.** If s,t are two characters of H, then the two representations  $\pi_s$  and  $\pi_t$  are unitarily equivalent if and only if their restrictions to the abelian subgroup H are unitarily equivalent. Denote by  $A \subset M$  the inclusion of von Neumann algebras  $L(H) \subset L(G)$ . The group normalizer  $N_G(H)$  generates the same von Neumann algebra than the normalizer of the algebras  $N_M(A)$ .

Proof. We recognize the Takesaki equivalence relation and the bimodule equivalence relation of  $L(H) \subset L(G)$  which are respectively  $\mathcal{R} = \{(s,t) \in \widehat{H}^2, \ \pi_s \simeq \pi_t\}$  and  $\mathcal{B} = \{(s,t) \in \widehat{H}^2, \ \pi_s|_{H} \simeq \pi_t|_{H}\}$ . Consider the orbit equivalence relation  $\mathcal{N}_0$  given by the action  $ad: \mathcal{N}_G(H) \curvearrowright \widehat{H}$ . If  $\mathcal{R} = \mathcal{B}$ , then we have the first assertion of the proposition. If  $\mathcal{N}_0 \equiv \mathcal{B}$ , then Proposition 2.2 implies the second assertion of the proposition. Let us show that  $\mathcal{R} = \mathcal{B} = \mathcal{N}_0$ . By definition,  $\mathcal{R} \subset \mathcal{B}$ . Let us show that  $\mathcal{B} \subset \mathcal{N}_0$ . Let  $(s,t) \in \mathcal{B}$ , and v a unitary of  $\ell^2(C)$  such that  $v^*\pi_s(h)v = \pi_t(h)$  for any  $h \in H$ . Consider  $h \in H$ ,

$$\pi_s(h)v(\varepsilon_1) = v\pi_t(h)(\varepsilon_1) = v(t(h)\varepsilon_1) = t(h)v(\varepsilon_1).$$

We denote by  $y = \sum_{c \in C} y_c \varepsilon_c$  the vector  $v(\varepsilon_1) \in \ell^2(C)$ , decomposed in the orthonormal basis  $\{\varepsilon_c, c \in C\}$ . Thus, for any  $h \in H$  and any  $c \in C$ ,

(3) 
$$y_c s(\eta(hc)) = y_{\sigma(hc)} t(h).$$

Let  $c \in C$  such that  $y_c \neq 0$ . The last identity tells us that for any h in H,  $|y_{\sigma(hc)}| = |y_c|$ . Therefore p(HcH) is finite, where  $p: G \twoheadrightarrow G/H$  is the canonical projection on the set of the right cosets. Condition (2) implies that c is in the normalizer  $N_G(H)$ . Equation (3) implies that  $s = ad_c(t)$ . Hence  $\mathcal{B} \subset \mathcal{N}_0$ .

Let us show that  $\mathcal{N}_0 \subset \mathcal{R}$ . Consider  $(s,t) \in \mathcal{N}_0$  and  $g \in \mathcal{N}_G(H)$  such that  $s = ad_g(t)$ . Let u be the unitary of  $\ell^2(C)$  defined as follows:

$$u(\varepsilon_c) = s(\eta(cg))\varepsilon_{\sigma(cg)},$$

where  $c \in C$ . An easy computation shows that  $u\pi_t u^* = \pi_s$ , hence  $\mathcal{N}_0 \subset \mathcal{R}$ .  $\square$ 

4.2. **Tensor product of MASAs.** In this section we compute the set of atoms for MASAs constructed from tensor product of MASAs. We deduce a result on the tensor product of von Neumann algebras generated by the normalizers.

Let  $\Lambda$  be a countable set and  $\{A_l \subset M_l, l \in \Lambda\}$  a family of MASAs in some finite von Neumann algebras. We fix a trace  $\tau_l$  on each  $M_l$ . Consider the infinite tensor products of von Neumann algebras with respect to the traces  $\tau_l$ 

$$\bigotimes_{l\in\Lambda}A_l\subset\bigotimes_{l\in\Lambda}M_l.$$

We denote those von Neumann algebras by  $A = \bigotimes_l A_l$  and by  $M = \bigotimes_l M_l$ . The Tomita commutant theorem implies that  $A \subset M$  is a MASA. For any  $l \in \Lambda$ , we consider a standard probability space  $(Y_l, \nu_l)$  and identify  $A_l$  with  $L^{\infty}(Y_l, \nu_l)$ . Let  $L^2(M_l)$  be the GNS Hilbert space and  $\pi_l, \rho_l : M_l \longrightarrow \mathbb{B}(L^2(M_l))$  the left and right actions of  $M_l$ . Consider the von Neumann algebra  $A_l = \{\pi_l(A_l) \cup \rho_l(A_l)\}'' \subset \mathbb{B}(L^2(M_l))$ . Let  $\mu_l$  be a measure on  $Y_l^2$  such that  $\mu_l(\Delta Y_l) = 1$  and such that  $A_l \simeq L^{\infty}(Y_l^2, \mu_l)$ , where  $\Delta Y_l$  is the diagonal. Consider a disintegration of  $\mu_l$  with respect to the projection  $p_l(s,t) = t$ ,

$$\mu_l = \int_{Y_l} \mu_{l,t} d\nu_l(t).$$

We denote the set of atoms of  $A_l \subset M_l$  by  $\mathcal{Y}_l$ . Let Y be the cartesian product of the  $Y_l$  with the  $\sigma$ -algebra generated by the subsets of Y of the form  $\prod_l X_l$ , where  $X_l \subset Y_l$  is a measurable subset that is equal to  $Y_l$  for all but a finite number of  $l \in \Lambda$ . Let  $\nu = \bigotimes_l \nu_l$  be the unique probability measure on Y that satisfies  $\nu(\prod_l X_l) = \prod_l \nu_l(X_l)$ . We identify A with  $L^{\infty}(Y, \nu)$ . We denote an element of Y by  $\underline{t}$  and its l-component by  $t_l$ .

**Theorem 4.2.** The set of atoms  $\mathcal{Y}$  of  $A \subset M$  is the set of couples  $(\underline{s},\underline{t})$  such that for any l  $(s_l,t_l) \in \mathcal{Y}_l$  and  $s_l = t_l$  for all but a finite number of  $l \in \Lambda$ .

Proof. Let  $\Omega^l$  be the image of the unity of  $M_l$  in the Hilbert space  $L^2(M_l)$ . Consider the infinite tensor product of Hilbert spaces  $\bigotimes_{l \in \Lambda} L^2(M_l)$  with respect to the vectors  $\Omega^l$ . There is a unitary transformation between  $L^2(M)$  and  $\bigotimes_{l \in \Lambda} L^2(M_l)$  that conjugates the actions  $\pi, \rho$  with the tensor product of actions  $\bigotimes_{l \in \Lambda} \pi_l$  and  $\bigotimes_{l \in \Lambda} \rho_l$ . For any  $l \in \Lambda$ , there exists a measurable field of Hilbert spaces  $\{\mathcal{H}^l_{s,t}, s, t \in Y_l\}$  such that the  $A_l$ -bimodule  $L^2(M_l)$  is isomorphic to the direct integral

$$\int_{Y_l^2}^{\oplus} \mathcal{H}_{s,t}^l d\mu_l(s,t).$$

Consider the disintegration of the vector  $\Omega^l$  which is

$$\int_{Y_l^2}^{\oplus} \Omega_{s,t}^l d\mu_l(s,t).$$

For any couple  $(\underline{s},\underline{t}) \in Y^2$ , we consider the infinite tensor product of Hilbert spaces  $\mathcal{H}_{\underline{s},\underline{t}} = \bigotimes_{l \in \Lambda} \mathcal{H}^l_{s_l,t_l}$  with respect to the vectors  $\Omega^l_{s_l,t_l}$ . For any finite subset  $E \subset \Lambda$ , there exists a unique measure  $\mu_E$  on  $Y^2$  that satisfies

$$\mu_E(\prod_{l\in\Lambda}X_l)=\prod_{l\in E}\mu_l(X_l)\times\prod_{l\in\Lambda\setminus E}\mu_l(X_l\cap\Delta Y_l).$$

Consider the measure class  $\mathcal{C}$  of  $Y^2$  generated by all the measures  $\mu_E$ . Let  $\mu$  be an element of the class  $\mathcal{C}$ . It is clear that the A-bimodule  $L^2(M)$  is isomorphic to the direct integral of Hilbert spaces  $\{\mathcal{H}_{\underline{s},\underline{t}}, \, (\underline{s},\underline{t}) \in Y^2\}$  over the measure  $\mu$ . It means that

$$L^2(M) \simeq \int_{Y^2}^{\oplus} \mathcal{H}_{\underline{s},\underline{t}} d\mu(\underline{s},\underline{t}).$$

Consider the projection of  $Y^2$   $p(\underline{s},\underline{t}) = \underline{t}$ . Let

$$\mu = \int_{Y} \mu_{\underline{t}} d\nu(\underline{t})$$

be a disintegration of  $\mu$  with respect to p and let  $\mathcal{Y}$  be the set of atoms of  $A \subset M$ . By the construction of  $\mu$ , it is clear that  $\mathcal{Y}$  is equal to the set of couples  $(\underline{s},\underline{t})$  such that for any l  $(s_l,t_l) \in \mathcal{Y}_l$  and  $s_l = t_l$  for all but a finite number of l.

**Corollary 4.3.** Let  $N_{M_l}(A_l)$  be the group normalizer of the MASA  $A_l \subset M_l$ , and consider the von Neumann algebra generated by those normalizers

$$\bigotimes_{l \in \Lambda} N_{M_l}(A_l)'' \subset \bigotimes_{l \in \Lambda} L^2(M_l).$$

Then,

$$N_M(A)^{\prime\prime} = \bigotimes_{l \in \Lambda} N_{M_l}(A_l)^{\prime\prime}.$$

In particular, a tensor product of singular MASAs is a singular MASA.

Proof. For any  $l \in \Lambda$  consider a countable subgroup  $G_l < N_{M_l}(A_l)$  such that  $G_l'' = N_{M_l}(A_l)'' \subset M_l$ . By Proposition 2.2 and Theorem 3.1 we have that the orbit equivalence relation  $\mathcal{N}_{G_l}$  is equivalent to  $\mathcal{Y}_l$ . Let G be the countable subgroup of  $N_M(A)$  generated by elements of the form  $u = \bigotimes_{l \in \Lambda} u_l$ , where  $u_l \in G_l$  and  $u_l = 1$  for all but finitely many  $l \in \Lambda$ . We see immediately that the orbit equivalence relation induced by G is equal to the set of couples  $(\underline{s},\underline{t})$  such that for any  $l \in \Lambda$ ,  $(s_l,t_l) \in \mathcal{N}_{G_l}$  and  $s_l = t_l$  for all but finitely many  $l \in \Lambda$ . Therefore Theorem 3.1 implies that  $\mathcal{N}_G = \mathcal{Y}$ . Hence, by Proposition 2.2 we have that  $G'' = N_M(A)'' \subset M$ . By construction

$$G'' = \bigotimes_{l \in \Lambda} G''_l = \bigotimes_{l \in \Lambda} N_{M_l}(A_l)'' \subset \bigotimes_{l \in \Lambda} L^2(M_l).$$

Therefore,

$$N_M(A)'' = \bigotimes_{l \in \Lambda} N_{M_l}(A_l)''.$$

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